Apparent Cosmic Acceleration via Backreaction from the Causal Propagation of Inhomogeneity Information

("Causal Backreaction")

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Department of Physics and Astronomy Hofstra University, NY, USA It's a Matter of Taste: Different Methods for Explaining the (Apparent?) Cosmic Acceleration Broadly Speaking, there are principal approaches:

# (1) Dark Energy:

 (i) Cosmological Constant (ΛCDM) → "Cosmological Constant Problem" & "Coincidence Problem"

(ii) Dynamical Dark Energy (DDE) → Yet Another Exotic Substance?
 (Nonadiabatic Pressure to stay smooth?)

- (2) Modified Gravity: f(R) Theories, etc.  $\rightarrow$  "Inelegant" modifications to G.R.?
- (3) **Inhomogeneities**: "Dressed" Cosmological Parameters -> Non-Copernican?
- (4) **Structure formation**: <u>*Backreaction*</u> on  $a(t)! \rightarrow$  Strong enough to work...???

Each approach has its own advantages... and its own problems...

After much subtle, sophisticated debate:

→ Everyone chooses their own favorite approach, anyway!

# Backreaction seems like the perfect solution!

... Automatically triggers at the right time, & with strength based on mass density  $\Omega_M$ ...

→ No coincidences! No changes to Einstein's G.R.! No special observers!

So why are the "naysayers" saying "<u>Nay</u>"?

Can the Acceleration of Our Universe Be Explained by the Effects of Inhomogeneities?

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#### Abstract

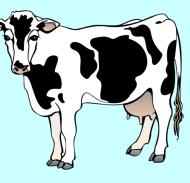
No. It is simply not plausible that cosmic acceleration could arise within the context of general relativity from a back-reaction effect of inhomogeneities in our universe, without the presence of a cosmological constant or "dark energy." We point out that our universe appears to be described very accurately on all scales by a Newtonianly perturbed FLRW metric. (This assertion is entirely consistent with the fact that we commonly encounter  $\delta \rho / \rho > 10^{30}$ .) If the universe is accurately described by a Newtonianly perturbed FLRW metric, then the back-reaction of inhomogeneities on the dynamics of the universe is negligible. If not, then it

• <u>The Central Issue</u>: In a Universe with (mostly) Nonrelativistic matter and (mostly) Newtonian Perturbations, how to get a *Strong-G.R.* effect like <u>Cosmic Acceleration</u>?

Formalisms/Models Lacking It\* <u>Necessary Physics</u> (as we'll see...) (\* Or, "How to offend every other researcher in Backreaction") **Overlapping/Cumulative** summing of Backreaction No-Go "Proofs" pert's. from different inhomogeneities (esp. from *outside* the "local matter horizon") Swiss-Cheese Models (Any Interior) **Vorticity** (and/or Velocity Dispersion) ...generated from... Purely Observational "Apparent" Structure Self-Stabilization & Virialization coll. (e.g., Voids, Lensing effects...) **Causal** Gravitational Info Propagation via: Terms at least up to  $O(v^2)$ Perturbation Theory Expansions **Tensor Components** (Most to date, If not all) "Magnetic" Gravitational Terms Buchert & Eblers formalism with Metric Pert. Potential Time-Derivatives "(Backreaction) as Total Divergence "Newtonian-Level" Strength Perturbations Static "Lattice of Inhomogeneities" calc's A Dynamical Phase Transition w/only the Final, Clumped "Steady State" from "Smooth" to "Clumped"

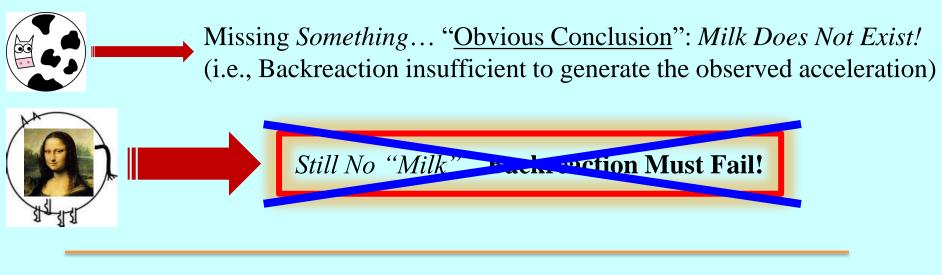
How Now Round Cow? : The Current State of Backreaction Research (as I see it ...)

Backreaction from Structure Formation, in the **Real Universe**:



The Physics of Backreaction, in most popular models:







"As far as the laws of mathematics refer to reality, they are not certain; and as far as they are certain, they do not refer to reality." - A. Einstein

Rather than doing exact models of very approximate physics, better to do a very approximate model of the exact physics!

# The <u>crucial difference</u> between *Instantaneous* "**Newtonian**(-Grav.) Cosmology", and *Newtonian-strength perturbations* in General Relativity, is often overlooked!

Important **caveat** ("the fine-print") from "<u>Averaging inhomogeneous Newtonian cosmologies</u>" (Buchert, T., & Ehlers, J. 1997, A&A 320, 1)

> An important difference to the Newtonian treatment, besides spatial curvature, arises due to the fact that it may not be in general possible to represent the term (18) as a divergence in GR. We stress that this would imply a strong challenge for the standard cosmologies, since we can no longer argue, except for non-generic situations, that there exist cases in which the average obeys Friedmann's law. Even more, we don't expect the previously discussed arguments (after eq. (18)) to hold, since the valid theory on the large scales under consideration is general relativity.

...and, from "<u>On average properties of inhomogeneous fluids in general relativity I:</u> <u>dust cosmologies</u>" (Buchert, T., 2000, Gen. Rel. Grav. 32, 105)

We conclude:

3. We were not able to produce an argument analoguous to the Newtonian treatment stating that the 'backreaction term' vanishes for topologically closed space sections, if integrated over the whole space. Without such an argument averaged inhomogeneous cosmologies cannot be identified with the standard FRW cosmologies on any spatial scale. To justify this identification as an approximation there is presently no sufficiently general quantitative result as to whether the 'backreaction' term could be neglected on some scale or, in words suggested by *Corollary 3*, whether the averaged curvature decouples from the inhomogeneities.

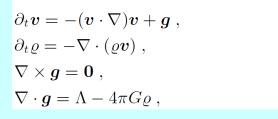
#### <u>Newtonian-Strength Perturbations</u>: Are they **Really** Negligible as a *Total Divergence*?

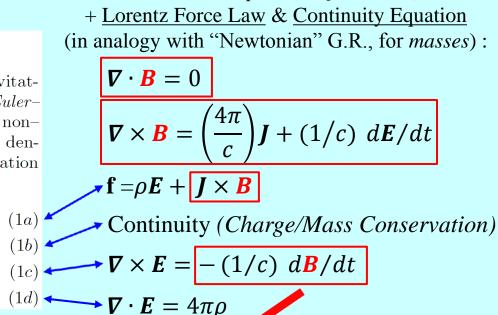
- Newtonian-level metric perturbations, apparently being expressible in the "Buchert formalism" as a total divergence, are believed to provide (essentially) *zero* Backreaction...
  - → Represents a huge impediment to acceleration-via-backreaction, since it implies the requirement of strongly non-Newtonian perturbations and relativistic flows!

From "<u>Averaging inhomogeneous Newtonian cosmologies</u>": (Buchert, T., & Ehlers, J. 1997, A&A 320, 1)

2. Averages in Newtonian cosmology

According to Newtonian physics, the motion of a self-gravitating, pressureless fluid ("dust") is governed by the *Euler-Poisson system* of equations. Thus, with respect to a nonrotating Eulerian coordinate system<sup>1</sup> the fields of mass density  $\rho(\boldsymbol{x},t) > 0$ , velocity  $\boldsymbol{v}(\boldsymbol{x},t)$  and gravitational acceleration  $\boldsymbol{g}(\boldsymbol{x},t)$  are required to satisfy





Maxwell's Equations (for E& M)

Something is <u>Missing</u>"... the "Magnetic" Gravitational Fields!

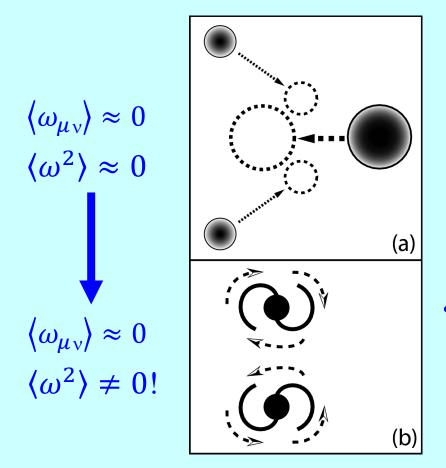
... No *B*-terms *>* No *Wave Propagation >* No *Gravitational Info* From <u>New</u>, *Distant* Structures!

"In the Newtonian approximation the expansion of a domain is influenced by the inhomogeneities inside the domain." (Buchert, Kerscher & Sicka, 2000, Phys. Rev. D62, 043525)

• "Causal Backreaction" is the idea that this view is unacceptable, even for Newtonian-Strength Pert's. !

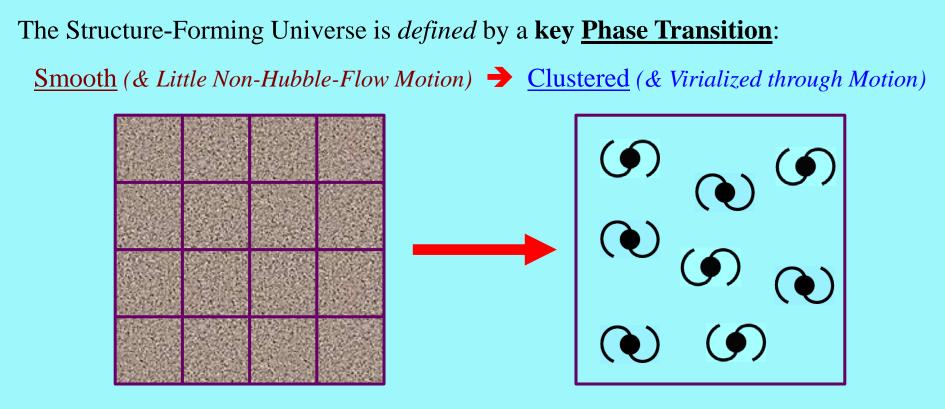
<u>Vorticity from Virialization</u>: The **Key Factor** vs. Gravity in Stabilizing All Structures

- Vorticity often *deliberately dropped* from calculations! (For convenience...?)
  - → Is Vorticity a "<u>Small-Scale Player</u>"? (e.g., Buchert, T. 2008, Gen. Rel. Grav. 40, 467) (*Relevant only for cosmic averages performed over domains* ≤ galaxy cluster scales??)
    - The quantity in the <u>Raychaudhuri Equation</u> for the increase of the velocity expansion ( $d\theta/dt$ ) is **vorticity squared** ( $\omega^2$ ), <u>not</u> the vorticity ( $\omega_{\mu\nu}$ )...



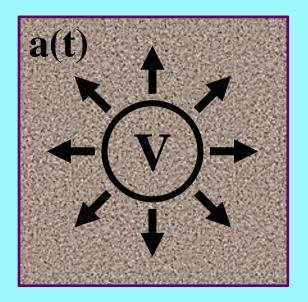
 $\rightarrow \langle \omega^2 \rangle$  obviously cannot average away!

- ...Cosmic averages of positive semi-definite quantities "attained in the subdomains is 'frozen' and cannot become smaller by averaging over larger domains." (Buchert & Ehlers, 1997, A&A 320, 1)
- Formalisms which neglect vorticity as a "small scale player" cannot properly estimate the effective backreaction from structure formation!



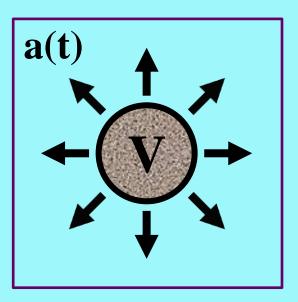
- Q: How best to model *all* of the relevant physics of this phase transition, without any **exact** (or "complete" perturbative) formalism that captures *everything*?
- <u>A</u>: Use the <u>relatively simple</u> nature of the **beginning** and **end** states of the structure-formation process to estimate the <u>net change</u> in the metric, **before**  $\rightarrow$  after.
  - The interim dynamics are less crucial, except to determine the *precise timing;* for now, will constrain this observationally as an *empirical* "<u>clumping function</u>".
  - Backreaction is a nonequilibrium process: it ends when the (causally-observed) structure formation is *complete*. ("Acceleration" just a comparison of "then" vs. "now"!)

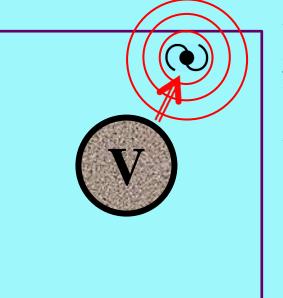
## Estimating the Net Effect of Clustering (on some "local volume", V):



As is well known for a <u>homogeneous</u> universe (e.g., Weinberg, 1972, "Gravitation and Cosmology"), the Friedmann expansion for V can be derived without reference to *anything* outside of it...

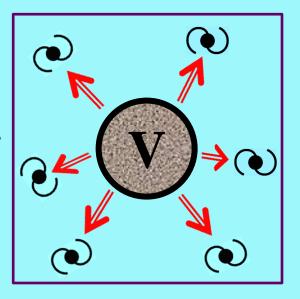
...so "**Remove**" the Exterior! : (<u>same</u> expansion behavior!)





When the universe becomes <u>inhomogeneous</u>, however, then *individually-clustered* & *vorticity-stabilized* objects become **gravitational attractors**, which pull on all other mass (*including* that within **V**)...

..and these grav. pulls upon **V** are <u>new</u>, as if the objects *"came in from infinity":* 



Therefore, can model the **main effect** (upon volume **V**) of <u>Fully-Virialized Clusters</u> by adding in the small, *Newtonian-strength* metric perturbation term for the mass of each self-stabilized system, *on top of* the internally-generated FRW metric for **V**:

 $g_{\mu\nu}(\mathbf{V}) = \{Unpert.FRW\} + \sum_{All\ Clumps,\ i} \{-dt^2 \left[-2GM_i(\mathbf{t})/a(t)r_i\right] + \frac{dr^2 a(t)^2 \left[2GM_i(\mathbf{t})/a(t)r_i\right]}{\mathbf{L}}\}$   $(Must\ be\ angle\ averaged\ for\ clumps\ in\ different\ directions})$ 

 $\rightarrow$  Each pert. term *slowly grows from zero* as each mass M<sub>i</sub> goes from **smooth** to **fully clumped**.

• <u>N.B.</u>: These perturbative factors (esp. the "<u>extra volume</u>" in  $g_{rr}$ ) are generated only because the clumps do not collapse completely, but *stabilize* themselves w/**vorticity**, velocity dispersion...

The gravitational *pulls* (forces) from clumps in *different directions* roughly cancel out in V;
 (a "<u>Smoothly-Inhomogeneous</u>" Universe); but the Potential Pert's., ΔΦ, *always add together!*

<u>Q</u>: Each individual pert. term is *very small*,  $\propto (1/r) \times \{ \text{orders } n > 1 \text{ of } (v/c)^n \}; \underline{can \text{ it matter}} \}$ ?

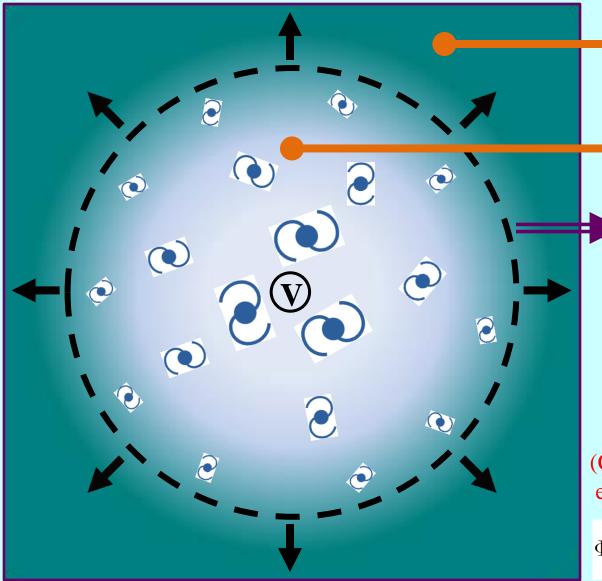
<u>A</u>: Yes, "little things" add up : (not like in <u>Swiss-Cheese</u> models!)

Tot. Pert. (at Obs.) 
$$\propto \int_{r=0}^{\infty} \frac{r^2 dr}{r} = \infty^2 !$$

→ A stronger divergence than **Olbers' Paradox**! (...only rendered finite by <u>causality</u>...)

Note, for all <u>Perturbation Theory</u> approaches: *"Small Amplitude"* terms **cannot** be reliably neglected, because for <u>cumulative</u> effects, size doesn't matter! <u>A Simple Analogy</u>: Gravitation is *much weaker*than Electromagnetism; is Gravity "<u>negligible</u>"? Is E&M holding you down in your seat right now?

#### But What Does an <u>Observer</u> See, Considering Causal "Look-Back Times"?



• Causal Backreaction effects are <u>finite</u>, though *large!* (& depend upon <u>faraway</u> pert's., not *local* clustering, as in Pert. Theory) Universe still *smooth*, **before** onset of clustering

• The later, clustered universe

"<u>Wave of Observed Clumpiness</u>" (the causal edge of clustering obs., *moving outward* at **c** ...)

→ "<u>Causal Backreaction</u>" is a relativistic process, even if most matter obeys v ≪ c !
 (Cannot drop 0[(v/c)<sup>2</sup>] terms or time derivatives!)

(Clustered mass density must be evaluated at *retarded time*)

$$\Phi_{\rm SR}(\mathbf{x},t) = -G \int^{\infty} \frac{[\rho(\mathbf{x}',t)]_{\rm ret}}{|\mathbf{x}-\mathbf{x}'|} d^3x'$$

"Causal Updating"

How to Implement Causal Backreaction: (B. Bochner: arXiv:1109.4686 & arXiv:1109.5155)

(1) Choose a "<u>Clumping Evolution Function</u>" to empirically model the time dynamics:

**<u>Try</u>:** (i) Linear Regime: " $\rho_{clump}$ " ~  $\delta\rho/\rho$  ~  $a(t) \propto t^{2/3}$ (ii) Nonlinear Regime: " $\rho_{clump}$ " ~  $\delta\rho/\rho \sim a(t)^{n\geq 3} \propto t^2$ (iii) Prop. to time for structures to form: " $\rho_{clump}$ " ~ t

(2) Compute the "Causal Updating" integral to get the total metric perturbation, *I(t)*, (at any location) in the past, as a fn. of time:

(3) Obtain the *final metric* for all <u>Cosmo Calc's</u>:  $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} [1 + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} + \{[a_{MD}(t)]^{2} + (1/3)I(t)]\} |d\vec{r}|^{2}$   $ds^{2} = -c^{2}[1 - I(t)] dt^{2} + \{[a_{MD}(t)]^{2} + (1/3)I(t)] dt^{2} + (1/3)I(t)] dt^{2}$ 

(4) Integrate the trajectory of a SN light-ray to calculate Luminosity Distance as a fn. of *z*, & thus get *many other* Cosmological Params.:

Test all models vs. SNe data, <u>optimizing</u>: a)  $t_{init} \equiv$  "Beginning" of Clustering b)  $\Psi_0 \equiv$  Clustered "Mass Fraction" Now

$$I(t) = \int_{0}^{\alpha_{\max}(t,t_{\min})} \{12 \ \Psi[t_{ret}(t,\alpha)] \ [(t_0/t)^{2/3}]\} \alpha \ d\alpha$$

$$(\alpha, t_{ret}) = (\alpha_{\max}, t_{\min})$$

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$$(\alpha, t_{ret})$$

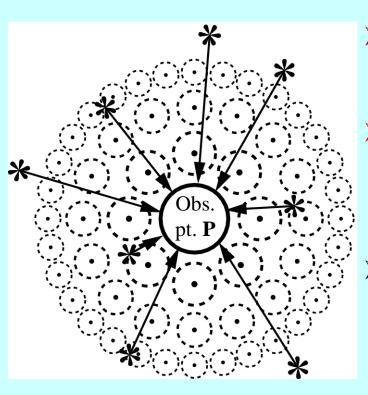
$$(r, t)$$

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$$(r, t)$$

**Important** Complication: *Old* metric perturbations from structures **slows down** all *new*, *ongoing* propagation of inhomogeneity information – <u>Weakening "Causal Updating"!</u>



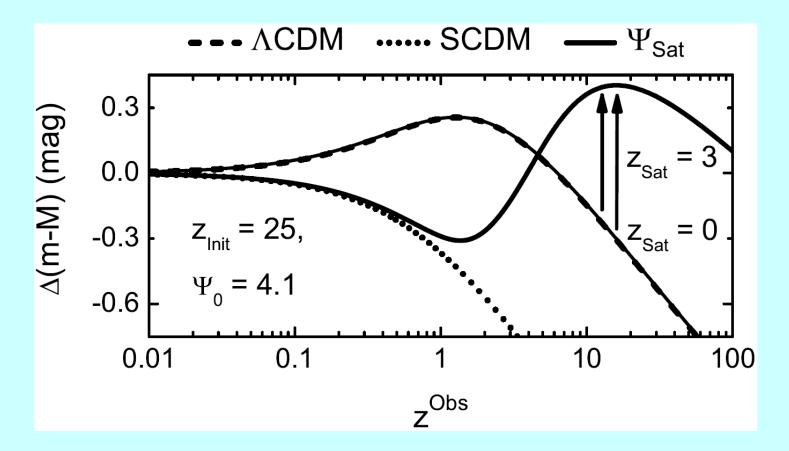
- Nothing other than Shapiro time delays on all propagating light & gravitational information.
- Accumulates over time... Causal Backreaction has a negative feedback loop upon itself...!
  - *"Eternal"* acceleration <u>not likely</u> here...
- This behavior is *recursive* later metric pert. effects depend upon prior ones, for a *nonlinear* response to clustering so call this "<u>Recursive Nonlinearities</u>" (to distinguish it from *nonlinear Gen. Rel.* effects).

*Version "<u>Cow 2.0</u>"*, now **including** Recursive Nonlinearities:

(B. Bochner: arXiv:1206.5056)

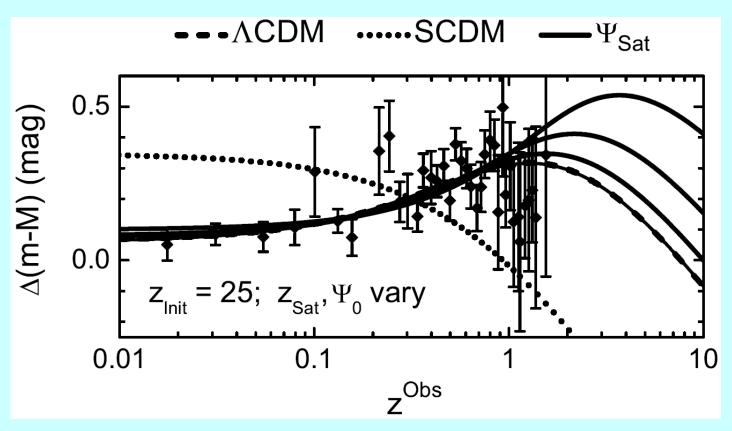
$$\begin{split} \alpha_{\max,i} &= \alpha_{\max,(i-1)} + \left\{ \frac{1}{3} \frac{\sqrt{1 - I_{(i-1)}^{\text{RNL}}}}{\sqrt{1 + [I_{(i-1)}^{\text{RNL}}/3]}} \frac{\Delta T}{[T_{(i-1)}]^{2/3}} \right\} ,\\ X_i^{\text{RNL}} &= \frac{12}{T_i^{2/3}} \sum_{k = \{1,(i-2)\}} \left\{ \Psi[T_{(i-k)}] [\alpha_{\max,i} - \alpha_{\max,(i-k)}] [\alpha_{\max,(i+1-k)} - \alpha_{\max,(i-k)}] \right\} \\ I_i^{\text{RNL}} &\sqrt{1 + [I_i^{\text{RNL}}/3]} = X_i^{\text{RNL}} . \end{split}$$

With Recursive Nonlinearities: Too Much Early Causal Backreaction Slows Down Later Effects!



- Causal Backreaction models **can** successfully mimic LCDM "Acceleration"! (...as long as structure formation & backreaction occur *gradually* enough...)
  - → Models from the Nonlinear Clustering Regime (% Clump  $\propto t^2$ ) are preferred
  - > But due to self-limiting feedback, **no** "Big Rip" likely from Causal Backreaction

Causal Backreaction models can reproduce the Apparent Acceleration (*Union Type Ia SNe*) for a "Smoothly-Inhomogeneous" Universe, without Dark Energy, Voids, etc.:



• Though successful at producing an "acceleration", need *large* clustering ( $\Psi_0 \sim 2 - 4$ )

- Equivalent to nearly-complete clustering on several different "hierarchical" scales simultaneously... (Stellar Clusters, Galaxies, Galaxy Clusters, etc... all individually stabilized through virialization)
- Less final clustering OK if using models where structure formation "<u>saturates</u>" late... (due to clumping-inhibiting feedback from "gastrophysics", and/or from the "acceleration" itself...)
  - Consistent with Vikhlinin, et al. (arXiv:0812.2720) results on recent effects on the growth of clustering

### Final Results for "Best-Fitting" models found (even without a rigorous optimization):

$z_{\mathrm{Sat}}$	$\Psi_{0,\mathrm{Opt}}$ <sup>a</sup>	$\chi^2_{ m Fit}$	$P_{\mathrm{Fit}}$ b	$I_0$	$z^{\mathrm{Obs}}$	$H_0^{\rm Obs}$	$H_0^{\rm FRW}$		$t_0^{\mathrm{Obs}}$	$\Omega_{\rm M}^{\rm FRW}$	$w_0^{\mathrm{Obs}}$	$j_0^{\rm Obs}$	$l_{\rm A}^{ m Obs}$	
				$\Psi_{\rm Sat}$ C	Clumping	Mode	l Runs,	$z_{]}$	$_{\rm init} = 2$	5				
0	4.1	311.8	0.351	0.53	1.14	70.07	42.32		13.64	0.943	-0.751	1.73	294.5	
0.25	2.6	313.5	0.326	0.58	1.15	69.60	40.24		14.00	1.054	-0.620	0.15	289.7	
0.5	2.3	316.6	0.284	0.68	1.15	69.40	36.32		14.65	1.338	-0.585	-0.14	279.8	
1	2.2	320.2	0.238	0.80	1.14	68.77	29.54		15.75	2.086	-0.488	-0.94	259.9	
	$Com_{l}$	parison	Values	from .	Best-Fit	flat $\Lambda$	CDM M	lι	$del \ (\Omega_{\ell})$	$\Lambda = 0.71$	3 = 1 - 9	$2_{\rm M})$		
•••	•••	311.9	0.380		1.0	69.96	69.96		13.64	0.287	-0.713	1.0	285.4	

Table 3: Output Cosmological Parameters from our RNL Runs with 'Early Saturation'

- Causal Backreaction models fit the (Union) SNe data essentially as well as flat  $\Lambda$ CDM
- A *sufficiently powerful* **Apparent Acceleration** now exists, to be consistent with important complementary Cosmological data sets
- There is a *significant difference* between the "<u>bare</u>" Hubble Constant  $(H_0^{FRW} \propto 1 / t_{FRW})$ , and the "<u>dressed</u>" Hubble Constant actually **observed**  $(H_0^{Obs})$ ; ... And so...
- Since the calculated "Critical Density" goes like  $\rho_{\text{Crit}} \propto H_0^2$ , and  $H_0^{\text{FRW}} \ll H_0^{\text{Obs}}$ , so that an *initially flat* <u>matter-only</u> universe ( $\Omega_M^{\text{FRW}} \equiv 1$ ) now **looks** under-dense ( $\Omega_M^{\text{Obs}} \sim 0.3$ ) !
- Other Cosmological Parameters also become Concordant without Dark Energy, such as the Observed Age of the Universe  $(t_0^{\text{Obs}})$  and the CMB 1<sup>st</sup> Peak Position  $(l_A^{\text{Obs}})$
- Deviations from  $\Lambda$ CDM can be tested via  $j_0^{Obs} \neq 1$  (though no iron-clad prediction yet)